

**Observation of the Production of Three Massive Gauge Bosons at  $\sqrt{s} = 13$  TeV**A. M. Sirunyan *et al.*\*  
(CMS Collaboration) (Received 19 June 2020; accepted 24 August 2020; published 5 October 2020)

The first observation is reported of the combined production of three massive gauge bosons ( $VVV$  with  $V = W, Z$ ) in proton-proton collisions at a center-of-mass energy of 13 TeV. The analysis is based on a data sample recorded by the CMS experiment at the CERN LHC corresponding to an integrated luminosity of  $137 \text{ fb}^{-1}$ . The searches for individual  $WWW$ ,  $WWZ$ ,  $WZZ$ , and  $ZZZ$  production are performed in final states with three, four, five, and six leptons (electrons or muons), or with two same-sign leptons plus one or two jets. The observed (expected) significance of the combined  $VVV$  production signal is 5.7 (5.9) standard deviations and the corresponding measured cross section relative to the standard model prediction is  $1.02^{+0.26}_{-0.23}$ . The significances of the individual  $WWW$  and  $WWZ$  production are 3.3 and 3.4 standard deviations, respectively. Measured production cross sections for the individual triboson processes are also reported.

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The production of three massive gauge bosons  $VVV$  ( $V = W, Z$ ) in high energy proton-proton ( $pp$ ) collisions is interesting because the standard model (SM) predictions for these processes involve the non-Abelian character of the theory [1]. In particular, the presence of quadruple gauge boson interactions can be probed through  $VVV$  production [2,3]. Triple gauge boson interactions and intermediate Higgs bosons ( $H$ ) also play a role. If physics beyond the SM is present at mass scales not far above 1 TeV, then cross section measurements for triple gauge boson production might deviate from SM predictions [4–7]. Up to now, such measurements have remained elusive because the production cross sections are low and backgrounds are insurmountable, except for rare leptonic final states. Next-to-leading order (NLO) SM calculations predict cross sections of 509, 354, 91.6, and 37.1 fb for  $WWW$ ,  $WWZ$ ,  $WZZ$ , and  $ZZZ$  production at 13 TeV with uncertainties of approximately 10% [8–12]. These calculations include contributions from the associated production of the Higgs boson with a  $V$  boson, where  $H$  decays to  $W^+W^-$  or  $ZZ$  [13–16]. In this analysis, these contributions generally are not the dominant ones, even though the cross section for  $VH$  production is relatively large, because the event selections described below have lower acceptances for the off-shell vector boson in  $H \rightarrow VV^*$ .

This Letter reports the first observation of  $VVV$  production in  $pp$  collisions at 13 TeV, using a data set

corresponding to an integrated luminosity of  $137 \text{ fb}^{-1}$ . Recently, the first evidence of  $VVV$  production in 13 TeV data was reported by the ATLAS Collaboration [17] following earlier searches for  $WWW$  production in 8 TeV ATLAS [18] and 13 TeV CMS data [19]. Five final states are considered (where  $\ell = e$  or  $\mu$ ):  $W^\pm W^\pm W^\mp \rightarrow \ell^\pm \ell^\pm 2\nu q\bar{q}'$ ,  $W^\pm W^\pm W^\mp \rightarrow \ell^\pm \ell^\pm \ell^\mp 3\nu$ ,  $W^\pm W^\mp Z \rightarrow \ell^\pm \ell^\mp 2\nu \ell^\pm \ell^\mp$ ,  $W^\pm ZZ \rightarrow \ell^\pm \nu 2(\ell^\pm \ell^\mp)$ , and  $ZZZ \rightarrow 3(\ell^\pm \ell^\mp)$ . This corresponds to five exclusive channels: two same-sign (SS) leptons with jets, three ( $3\ell$ ), four ( $4\ell$ ), five ( $5\ell$ ), and six ( $6\ell$ ) leptons. Searches in the dilepton and trilepton final states target  $WWW$  production; four-lepton events are used to search for  $WWZ$  production; and five- and six-lepton events are used to search for  $WZZ$  and  $ZZZ$  production, respectively.

The data were recorded in 2016–2018 with the CMS detector, whose central feature is a superconducting solenoid of 6 m internal diameter providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two end cap sections. Forward calorimeters extend the pseudorapidity ( $\eta$ ) coverage provided by the barrel and end cap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid. Events are selected using triggers [20] that require two electrons, two muons, or one electron and one muon passing loose isolation requirements and certain transverse momentum ( $p_T$ ) thresholds. A detailed description of the detector and definitions of the coordinate system are given in Ref. [21].

The CMS event reconstruction is based on the particle-flow (PF) algorithm [22], which combines information from the tracker, calorimeters, and muon systems to identify charged and neutral hadrons, photons, electrons, and muons,

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known collectively as PF candidates. Electrons and muons from  $V$  decays, known as prompt leptons, are selected for off-line analysis using standard criteria [23,24]. Events containing  $\tau$  leptons decaying into charged hadrons are rejected by requiring the absence of isolated tracks aside from selected electrons and muons. The PF candidates are clustered into jets using the anti- $k_T$  algorithm with a distance parameter of 0.4 [25–27]. Jets with  $p_T > 20$  GeV and  $|\eta| < 5$  are selected for the analysis. Defining the distance between a jet and a selected lepton by  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$  where  $\phi$  is the azimuthal angle, jets are rejected if  $\Delta R < 0.4$ . Jets containing the decay of a  $b$  quark are identified using the loose working point of the deep combined secondary-vertex  $b$  tagging algorithm [28]. To increase the efficiency for identifying low- $p_T$   $b$  hadrons not clustered into jets, a soft  $b$  tag object [29] is defined using a track-based secondary vertex reconstruction.

The primary  $pp$  interaction vertex is the reconstructed vertex with the largest summed  $p_T^2$  calculated using track-based jets and the associated missing transverse momentum ( $\vec{p}_T^{\text{miss}}$ ), the negative  $\vec{p}_T$  sum of those jets [30]. Track-based jets are constructed using only tracks associated with the given vertex. In addition to the primary interaction, other  $pp$  interactions (pileup) produce extra charged particles and neutral energy. Only tracks associated with the primary vertex are used. The average neutral energy density from pileup is estimated, and subtracted from the reconstructed jet energies and the energy sum used in the calculation of lepton isolation [31].

The previous search for  $WWW$  production [19] is based on sequences of requirements called sequential cuts. In this Letter, that approach is extended to cover all five channels. In addition, motivated by the relatively high yields in the SS,  $3\ell$ , and  $4\ell$  channels, multivariate techniques based on boosted decision trees (BDTs) [32–36] are applied that outperform the sequential-cut analyses. Both the sequential-cut and BDT-based analyses are presented.

The acceptances, efficiencies, and kinematic properties of the signal and background processes are determined using a combination of data and simulated events. The POWHEG 2.0 [37–40] and the MadGraph5\_aMC@NLO (2.2.2 and 2.4.2) generators [41] are used to generate  $VVV$  signal events (including  $VH$ ), diboson ( $VV$ ), and single- $t$  background events. The MadGraph5\_aMC@NLO generator is used in the leading-order (LO) mode with MLM jet matching [42] to generate SM  $t\bar{t}$ ,  $t\bar{t} + X$  ( $X = W, Z, H$ ),  $W + \text{jets}$ ,  $Z + \text{jets}$ ,  $W\gamma$ , and  $W^\pm W^\pm$  events. The most precise cross section calculations available are used to normalize the simulated samples, and usually correspond to either NLO or next-to-NLO accuracy [16,41,43–50]. Parton showering, hadronization, and the underlying event are modeled by PYTHIA (8.205 and 8.230) [51] with parameters set by the CUETP8M1 [52] and CP5 tune [53]. The NNPDF 3.0 [54] and 3.1 [55] parton distribution functions (PDFs) are used in the generation of all simulated samples. Pileup is

simulated and the GEANT4 [56] package is used to mimic the response of the CMS detector.

The SS channel targets  $WWW$  production [19] and requires exactly two SS leptons with  $p_T > 25$  GeV and one or more jets. The dilepton mass  $m_{\ell\ell}$  must exceed 20 GeV. This channel is subdivided into nine signal regions according to the flavors of the leptons ( $e^\pm e^\pm$ ,  $e^\pm \mu^\pm$ , or  $\mu^\pm \mu^\pm$ ) and the jet content. Events with exactly one jet are denoted “1 $J$ .” Events with two or more jets are categorized as “ $m_{jj}$  -in” or “ $m_{jj}$  -out” depending on whether the dijet mass for the two jets closest in  $\Delta R$  is compatible with the  $W$  boson mass ( $65 < m_{jj} < 95$  GeV). The background processes fall broadly into three categories. The first category contains trilepton processes with one lepton either not selected or not reconstructed (“lost”). Such backgrounds include  $WZ$  and  $t\bar{t}Z$  production, which typically have only one prompt neutrino in the final state; they are reduced by requiring  $m_T^{\text{max}} > 90$  GeV, where  $m_T^{\text{max}}$  is the largest transverse mass obtained from  $\vec{p}_T^{\text{miss}}$  and any lepton in the event. The second category consists of processes with SS lepton pairs, mainly from  $W^\pm W^\pm + \text{jets}$  and  $t\bar{t}W^\pm$  production. This contribution is suppressed by requiring  $m_{JJ} < 500$  GeV and  $|\Delta\eta_{JJ}| < 2.5$  for the two highest- $p_T$  jets. The third category includes  $W + \text{jets}$  and  $t\bar{t} + \text{jets}$  production where a final-state jet or photon is misidentified as a charged lepton, and is labeled nonprompt. These background contributions are suppressed using strict lepton identification and isolation requirements and by requiring  $p_T^{\text{miss}} > 45$  GeV. All backgrounds containing top quarks are further reduced by excluding events with  $b$ -tagged jets or soft  $b$  tags. The background due to charge mismeasurement in Drell-Yan production is relevant only for dielectron events and is reduced to a negligible level by requiring  $|m_{\ell\ell} - m_Z| > 10$  GeV.

The  $3\ell$  channel, which also targets  $WWW$  production, is subdivided according to the number of same-flavor opposite-sign (SFOS) lepton pairs: 0SFOS, 1SFOS, and 2SFOS. At least one lepton is required to have  $p_T > 25$  GeV, while the others must have  $p_T > 20$  GeV, except in 0SFOS where all three leptons are required to have  $p_T > 25$  GeV to reduce contamination from nonprompt leptons. Events in 1SFOS and 2SFOS must contain no jets, whereas the presence of one jet is allowed in 0SFOS. The background sources are similar to those in the SS category. Events with  $b$ -tagged jets are excluded to suppress the nonprompt-lepton background from processes involving top quarks. The contribution from triple prompt lepton backgrounds is suppressed by requiring  $|m_{\ell\ell} - m_Z| > 20$  GeV and  $m_{\ell\ell} > 20$  GeV for all SFOS pairs. Additional background reduction is achieved with the following requirements: if exactly one SFOS lepton pair is found then  $m_T^{\text{3rd}}$ , defined as the transverse mass calculated from  $\vec{p}_T^{\text{miss}}$  and the third lepton that is not one of the SFOS pair, must be larger than 90 GeV; and, for events with no SFOS pairs,  $m_T^{\text{max}} > 90$  GeV is required. Background contributions from

nonprompt leptons and converted or misidentified photons are reduced by requiring a large  $p_T$  of the three-lepton system  $|\vec{p}_T^{3\ell}| > 50$  GeV, and a large azimuthal separation  $\Delta\phi(\vec{p}_T^{3\ell}, \vec{p}_T^{\text{miss}}) > 2.5$  between  $\vec{p}_T^{\text{miss}}$  and  $\vec{p}_T^{3\ell}$ . Events with a conversion photon emitted in a  $Z$  boson decay are suppressed by requiring  $|m_{3\ell} - m_Z| > 10$  GeV where  $m_{3\ell}$  is the three-lepton invariant mass.

The  $4\ell$  channel targets  $WWZ$  production. The  $Z$  boson is identified through its decay to a SFOS lepton pair with  $|m_{\ell\ell} - m_Z| < 10$  GeV. These leptons are required to have  $p_T > 25(10)$  GeV for the (sub)leading lepton. The (sub)leading lepton of the remaining non- $Z$  leptons must have  $p_T > 25(10)$  GeV. The dominant background comes from  $ZZ$  production, so the cases of different-flavor ( $e\mu$ ) and same-flavor ( $ee/\mu\mu$ ) non- $Z$  lepton pairs are handled separately. The non- $Z$  same-flavor invariant mass is required to differ from  $m_Z$  by at least 10 GeV. Other background contributions consist of  $t\bar{t}Z$ ,  $tWZ$ ,  $t\bar{t}H$ , and  $WZ$  events. The rejection of events with  $b$ -tagged jets reduces contributions from top quarks and a requirement that  $m_{\ell\ell} > 12$  GeV for all opposite-sign lepton pairs suppresses backgrounds from low-mass resonances. The  $4\ell$  channel is subdivided into seven signal regions: for the  $e\mu$  category there are four bins in  $m_{\ell\ell}$  and  $m_{T2}$  [57], and for the  $ee/\mu\mu$  category there are three bins based on  $p_T^{4\ell}$  and  $p_T^{\text{miss}}$ .

The  $5\ell$  and  $6\ell$  channels target  $WZZ$  and  $ZZZ$  production, respectively. Event yields are low because of small cross sections and branching fractions. Since background contributions are low, the selection maximizes the signal efficiency. The two leading leptons are required to have  $p_T > 25$  GeV and other leptons must have  $p_T > 10$  GeV. Events in the  $5\ell$  channel are required to contain two SFOS lepton pairs with  $|m_{\ell\ell} - m_Z| < 15$  GeV. The background in the  $5\ell$  channel consists almost entirely of  $ZZ$  events with a nonprompt lepton, which is usually an electron. The background is reduced by requiring  $m_T > 50$  GeV, where  $m_T$  is calculated from  $\vec{p}_T^{\text{miss}}$  and that electron. Smaller background contributions arise from  $t\bar{t}Z$  and  $t\bar{t}H$  production, which are reduced by rejecting events with  $b$ -tagged jets. Events in the  $6\ell$  channel are required to have three SFOS pairs and a six-lepton scalar  $p_T$  sum larger than 250 GeV. The small  $6\ell$  background comes from  $t\bar{t}H$  and  $ZZ$  production.

Background contributions from sources with a particular number of prompt leptons and no nonprompt leptons in signal regions are estimated using simulations with correction factors, typically near unity, derived from several control regions in data enriched in the main sources of background events. Both the predicted numbers of events and relevant kinematic distributions are compared with observations in control regions to derive the correction factors. The precision of the comparison is used to assess systematic uncertainties in these background contributions. Background contributions from sources with one or more nonprompt leptons cannot be reliably evaluated using simulations, so estimates based on control samples in data

are used instead. These estimates rely on the fact that nonprompt leptons tend to be less isolated than prompt ones. For the SS and  $3\ell$  channels, following Ref. [19], the contribution of events with a nonprompt lepton is evaluated using a sample of events in which one lepton satisfies loose identification criteria but fails the tight criteria. The number of events in this region determines the estimate of the nonprompt background in the signal region using a transfer factor computed with a separate event sample rich in nonprompt leptons. This transfer factor is the ratio of the number of events that pass the tight selection criteria to those that pass the loose criteria. For the  $5\ell$  channel, a sample of events with three prompt leptons and one nonprompt lepton is dominated by  $WZ$  production and used to verify the prediction of background contributions with nonprompt leptons. Nonprompt leptons are a minor background for all other channels.

The signal strength  $\mu$ , defined as the measured production cross section times branching fraction divided by the expected SM value, is determined through simultaneous fits to all twenty-one signal regions. In one version of the fit, four independent signal strengths ( $\mu_{WWW}$ ,  $\mu_{WWZ}$ ,  $\mu_{WZZ}$ , and  $\mu_{ZZZ}$ ) are used. In the other version, a common signal strength  $\mu_{\text{comb}}$  is used for all four processes.

The most important sources of systematic uncertainty involve the estimation of background contributions; the uncertainties range from 5% to 25% and come mainly from limited statistical precision in the control regions. The uncertainties in the nonprompt background estimates from control samples in data also contribute significantly at 50%. Uncertainties related to trigger efficiencies, lepton identification and energy resolution, jet energy scale, and  $b$ -jet tagging efficiency range from 1% to 9%. A 2.3%–2.5% uncertainty in the integrated luminosity is assessed [58–60]. Uncertainties due to limitations of the theory include missing higher-order corrections (2%–14%), PDF uncertainties (2%–7%), and the strong coupling  $\alpha_s$  (1%). Theoretical and experimental uncertainties are correlated across different channels. Statistical uncertainties are much larger than systematic ones. The expected significance of the combined  $VVV$  production signal based on the sequential-cut selection is 5.4 standard deviations (s.d.), and the observed significance is 5.0 s.d. The observed (expected) significances for the individual triboson production processes are 2.5 (2.9) s.d. for  $WWW$ , 3.5 (3.6) s.d. for  $WWZ$ , 1.6 (0.7) s.d. for  $WZZ$ , and 0.0 (0.9) s.d. for  $ZZZ$ .

The discrimination of signal and background events in the SS,  $3\ell$ , and  $4\ell$  channels is enhanced by using BDTs. The training and optimization of the BDTs is carried out for each channel using simulated background and signal events. A minimum value of each BDT output variable substitutes for the categorizations of events and the kinematic requirements applied in the sequential-cut analyses. In the SS and  $3\ell$  channels, two separate BDTs are trained: the first one to separate the signal from the nonprompt

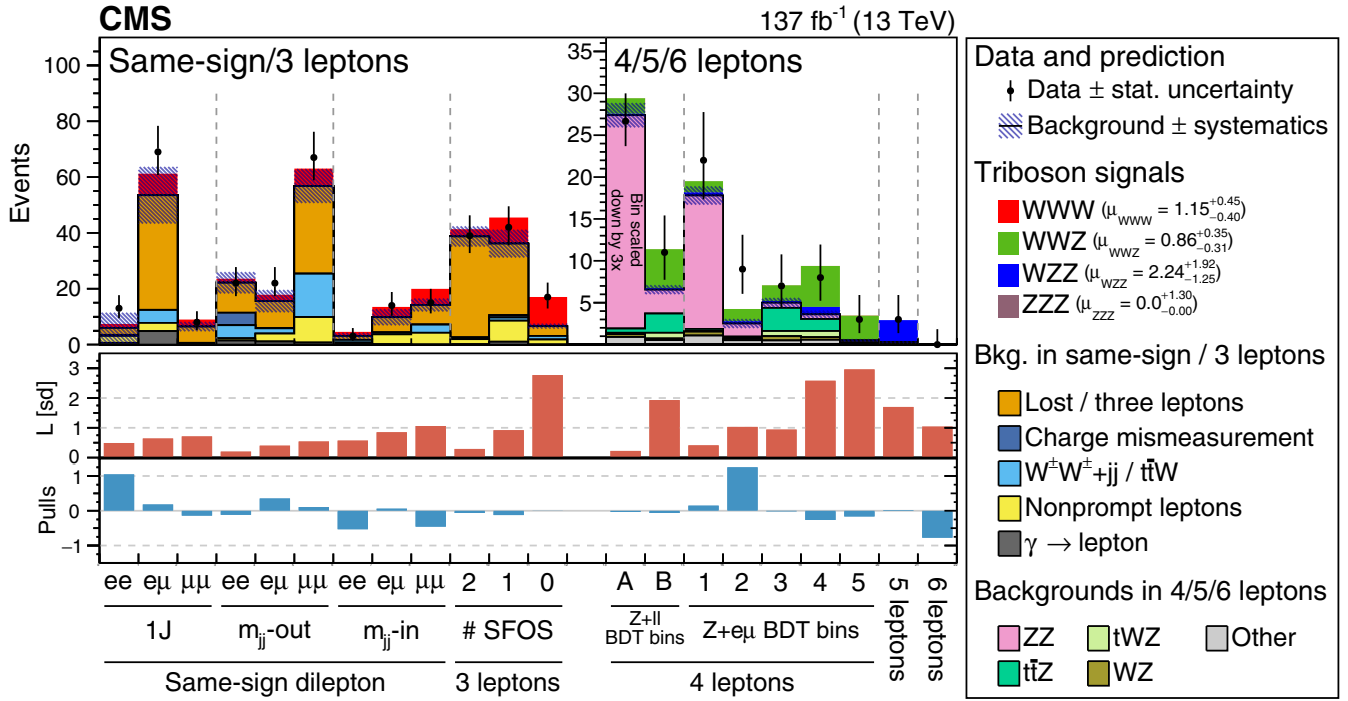


FIG. 1. Comparison of the observed numbers of events to the predicted yields after fitting. For the  $WWW$  and  $WWZ$  channels, the results from the BDT-based selections are used. The  $VVV$  signal is shown stacked on top of the total background. The points represent the data and the error bars show the statistical uncertainties. The expected significance  $L$  in the middle panel represents the number of standard deviations (s.d.) with which the null hypothesis (no signal) is rejected; it is calculated for the fit for  $\mu_{\text{comb}}$ . The lower panel shows the pulls for the fit result.

background and the second one to separate the signal from the rest of the background. These two BDTs are applied sequentially. In the  $4\ell$  channel, a similar strategy is pursued except that the two BDTs are targeted against  $ZZ$  and  $t\bar{t}Z$  backgrounds specifically. There are two (five) signal regions for events in the  $ee/\mu\mu$  ( $e\mu$ ) category. The improvement in sensitivity due to the use of BDTs varies channel by channel and is in the range 5%–15%. No BDTs are used for the  $5\ell$  and  $6\ell$  channels.

The yields in the individual signal regions obtained using the BDTs are shown in Fig. 1. The significances  $L$  of the expected numbers of events are computed including systematic uncertainties and are evaluated under the asymptotic approximation [61]. Pulls are the differences in the numbers of observed and predicted events normalized to the uncertainties in the numbers of predicted events. Assuming the SM production of  $VVV$  events, the expected significance of the fit with a single signal strength  $\mu_{\text{comb}}$  is 5.9 s.d. and the observed significance is 5.7 s.d. The observed (expected) significances for the individual triboson production processes are 3.3 (3.1) s.d. for  $WWW$ , 3.4 (4.1) s.d. for  $WWZ$ , 1.7 (0.7) s.d. for  $WZZ$ , and 0.0 (0.9) s.d. for  $ZZZ$ . In the most sensitive signal regions, approximately one third of the  $VVV$  events come from  $VH$  production. The measured signal strengths, obtained in the asymptotic approximation of the  $CL_s$  method [61], correspond to the total cross sections listed in Table I; leptonic

branching fractions for  $W$  and  $Z$  decays come from Ref. [62]. If  $VH$  is considered as a background, then the combined observed (expected) significance for  $\mu_{\text{comb}}$  is 2.9 (3.5) s.d. and the measured cross sections are listed in Table I. For  $ZZZ$  production, upper limits are reported at 95% confidence level. Signal strengths obtained using both sequential-cut and BDT-based approaches and with  $VH$  production counted as signal are summarized in Fig. 2.

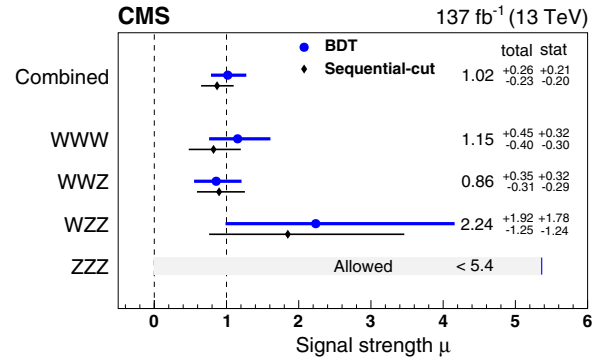


FIG. 2. Best fit values of the signal strengths for the BDT-based analyses (blue solid circles) and the sequential-cut analyses (black solid diamonds). The error bars represent the total uncertainty. For  $ZZZ$  production, a 95% confidence level upper limit is shown. The stated numerical values correspond to the BDT-based analysis.

TABLE I. Measured cross sections obtained with the BDT-based analyses. The uncertainties listed are statistical and systematic. For the results listed in the middle (right) column of the table, Higgs boson contributions are counted as signal (background). The  $VVV$  cross section is calculated from the fit for  $\mu_{\text{comb}}$ . For  $ZZZ$  production, 95% confidence level upper limits are reported.

Process	Cross section (fb)	
	Treating Higgs boson contributions as	
	Signal	Background
$VVV$	$1010^{+210+150}_{-200-120}$	$370^{+140+80}_{-130-60}$
$WWW$	$590^{+160+160}_{-150-130}$	$190^{+110+80}_{-100-70}$
$WWZ$	$300^{+120+50}_{-100-40}$	$100^{+80+30}_{-70-30}$
$WZZ$	$200^{+160+70}_{-110-20}$	$110^{+100+30}_{-70-10}$
$ZZZ$	$< 200$	$< 80$

In summary, proton-proton collision data at  $\sqrt{s} = 13$  TeV recorded with the CMS experiment and amounting to  $137 \text{ fb}^{-1}$  were used to observe the production of three massive gauge bosons. The significance of the observation is 5.7 standard deviations (s.d.) with 5.9 s.d. expected. For  $WWW$  ( $WWZ$ ) production, the observed significance is 3.3 (3.4) s.d. compatible with 3.1 (4.1) s.d. expected. Measured cross sections for  $WWW$ ,  $WWZ$ , and  $WZZ$  production and an upper limit for  $ZZZ$  production are in agreement with the expectations of the standard model. This Letter documents the evidence for  $WWW$  and  $WWZ$  production and the first observation of the combined production of three massive gauge bosons.

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 H. Kwon,<sup>90</sup> H. Lee,<sup>90</sup> K. Lee,<sup>90</sup> S. Lee,<sup>90</sup> K. Nam,<sup>90</sup> B. H. Oh,<sup>90</sup> M. Oh,<sup>90</sup> S. B. Oh,<sup>90</sup> B. C. Radburn-Smith,<sup>90</sup> H. Seo,<sup>90</sup>  
 U.K. Yang,<sup>90</sup> I. Yoon,<sup>90</sup> D. Jeon,<sup>91</sup> J. H. Kim,<sup>91</sup> B. Ko,<sup>91</sup> J. S. H. Lee,<sup>91</sup> I. C. Park,<sup>91</sup> Y. Roh,<sup>91</sup> D. Song,<sup>91</sup> I. J. Watson,<sup>91</sup>  
 H. D. Yoo,<sup>92</sup> Y. Choi,<sup>93</sup> C. Hwang,<sup>93</sup> Y. Jeong,<sup>93</sup> H. Lee,<sup>93</sup> J. Lee,<sup>93</sup> Y. Lee,<sup>93</sup> I. Yu,<sup>93</sup> V. Veckalns,<sup>94,pp</sup> A. Juodagalvis,<sup>95</sup>  
 A. Rinkevicius,<sup>95</sup> G. Tamulaitis,<sup>95</sup> W. A. T. Wan Abdullah,<sup>96</sup> M. N. Yusli,<sup>96</sup> Z. Zolkapli,<sup>96</sup> J. F. Benitez,<sup>97</sup>  
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 M. Ramirez-Garcia,<sup>99</sup> F. Vazquez Valencia,<sup>99</sup> J. Eysermans,<sup>100</sup> I. Pedraza,<sup>100</sup> H. A. Salazar Ibarguen,<sup>100</sup> C. Uribe Estrada,<sup>100</sup>  
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 M. I. Asghar,<sup>105</sup> M. I. M. Awan,<sup>105</sup> Q. Hassan,<sup>105</sup> H. R. Hoorani,<sup>105</sup> W. A. Khan,<sup>105</sup> M. A. Shah,<sup>105</sup> M. Shoaib,<sup>105</sup>  
 M. Waqas,<sup>105</sup> V. Avati,<sup>106</sup> L. Grzanka,<sup>106</sup> M. Malawski,<sup>106</sup> H. Bialkowska,<sup>107</sup> M. Bluj,<sup>107</sup> B. Boimska,<sup>107</sup> T. Frueboes,<sup>107</sup>  
 M. Górski,<sup>107</sup> M. Kazana,<sup>107</sup> M. Szeleper,<sup>107</sup> P. Traczyk,<sup>107</sup> P. Zalewski,<sup>107</sup> K. Bunkowski,<sup>108</sup> A. Byszuk,<sup>108,rr</sup> K. Doroba,<sup>108</sup>  
 A. Kalinowski,<sup>108</sup> M. Konecki,<sup>108</sup> J. Krolikowski,<sup>108</sup> M. Olszewski,<sup>108</sup> M. Walczak,<sup>108</sup> M. Araujo,<sup>109</sup> P. Bargassa,<sup>109</sup>  
 D. Bastos,<sup>109</sup> A. Di Francesco,<sup>109</sup> P. Faccioli,<sup>109</sup> B. Galinhas,<sup>109</sup> M. Gallinaro,<sup>109</sup> J. Hollar,<sup>109</sup> N. Leonardo,<sup>109</sup> T. Niknejad,<sup>109</sup>  
 J. Seixas,<sup>109</sup> K. Shchelina,<sup>109</sup> O. Toldaiev,<sup>109</sup> J. Varela,<sup>109</sup> S. Afanasiev,<sup>110</sup> P. Bunin,<sup>110</sup> M. Gavrilenko,<sup>110</sup> I. Golutvin,<sup>110</sup>  
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 V. Palichik,<sup>110</sup> V. Perelygin,<sup>110</sup> M. Savina,<sup>110</sup> D. Seitova,<sup>110</sup> V. Shalaev,<sup>110</sup> S. Shmatov,<sup>110</sup> S. Shulha,<sup>110</sup> V. Smirnov,<sup>110</sup>  
 O. Teryaev,<sup>110</sup> N. Voytishin,<sup>110</sup> A. Zarubin,<sup>110</sup> I. Zhizhin,<sup>110</sup> G. Gavrillov,<sup>111</sup> V. Golovtsov,<sup>111</sup> Y. Ivanov,<sup>111</sup> V. Kim,<sup>111,uu</sup>  
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 M. Toms,<sup>113</sup> E. Vlasov,<sup>113</sup> A. Zhokin,<sup>113</sup> T. Aushev,<sup>114</sup> R. Chistov,<sup>115,xx</sup> M. Danilov,<sup>115,xx</sup> A. Oskin,<sup>115</sup> P. Parygin,<sup>115</sup>  
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 MK Jayananda,<sup>127</sup> B. Kailasapathy,<sup>127,ccc</sup> D. U. J. Sonnadara,<sup>127</sup> DDC Wickramaratna,<sup>127</sup> W. G. D. Dharmaratna,<sup>128</sup>  
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